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Direct observation of structural changes near and
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The detailed structural changes that occur during and following the formation of a long-range, charge-ordered state have been accurately and rapidly determined. The localisation of the electron in ilvaite, $\text{Ca}(\text{Fe,Mn})_3\text{Si}_2\text{O}_8(\text{OH})$, forming Fe^{2+} and Fe^{3+} from the delocalized $2 \times \text{Fe}^{2.5+}$, is directly apparent in bond valence sums derived from accurately determined Fe–O distances extracted from neutron powder diffraction data obtained for 40 temperatures around the charge ordering temperature in just a few hours. Subtle changes that occur in the structure and lattice parameters after the formation of the charge-ordered state are also apparent for the first time from this work.

Charge-ordering behaviours occurring in materials represent an area of intense current research interest including studies of systems such as the manganites, LaMnO_3 , $\text{Pr}_{0.5}\text{Sr}_{0.5}\text{MnO}_3$ and $\text{La}_{1.2}\text{Sr}_{1.8}\text{Mn}_2\text{O}_7$,¹ CaFeO_3 ,² $\text{TbBaFe}_2\text{O}_5$,³ Fe_2OBO_3 ,⁴ LuFe_2O_4 ,⁵ $\text{Ca}_{2-x}\text{Na}_x\text{CuO}_2\text{Cl}_2$,⁶ Fe_3O_4 ,⁷ $\text{La}_{1.67}\text{Sr}_{0.33}\text{NiO}_4$,⁸ and YBaCo_2O_5 .⁹ In such systems the direct investigation of the charge-ordering transition is best achieved through the application of structure determination techniques, such as X-ray diffraction, neutron diffraction and EXAFS, which all show the gross changes in structure that occur as the charge localizes. Of these techniques neutron diffraction provides the most accurate information as the charge-ordering phenomenon is normally associated with small changes in metal to oxygen bond lengths as a result of electron localization; analysis of neutron diffraction data produces accurate and precise definition of the oxygen nuclear positions in the presence of heavier metals, derived from the large scattering length of oxygen. While single crystal neutron diffraction data provide the most accurate structure models, the requirement for large crystals and very long experimental times, where a significant temperature range has to be studied, often

prohibits such experiments. Profile fitting of neutron powder diffraction (NPD) data also produces accurate metal–oxygen distances but low neutron fluxes require extended experimental count times, typically a few hours per full diffraction profile; this has again limited the range and resolution of variable temperature work covering a full charge-ordering event so, for example, published NPD work on YBaCo_2O_5 ,⁹ and Fe_2OBO_3 ,⁴ were both limited to just a few temperatures though the recent study of short-range charge ordering in $\text{La}_{1.67}\text{Sr}_{0.33}\text{NiO}_4$,⁸ completed over 3 days on a 15 g sample is noteworthy. Careful and detailed studies of charge-ordering in materials, such as with $\text{TbBaFe}_2\text{O}_5$,³ can show additional anomalies in physical properties and unit cell volumes associated with premonitory or post-monitory transitions derived from local order or relaxation – however, structural changes associated with such transitions have not been resolved before. In this communication we report a detailed study of the charge ordering behaviour in $\text{CaFe}_2(\text{Fe,Mn})\text{Si}_2\text{O}_8(\text{OH})$, which occurs as the natural mineral ilvaite, using a very high flux NPD instrument on a 1.5 g sample in a 6 h experiment. The use of a high flux NPD instrument allows an unprecedented accurate resolution of the charge-ordering process enabling direct observation of the electron localization through determination of site changes and, for the first time, investigation of the subtle changes that occur in the structure just below the charge-ordering temperature.

Ilvaite, $\text{CaFe}_2\text{FeSi}_2\text{O}_8(\text{OH})$, is a mixed valence Fe-bearing naturally-occurring sorosilicate, in which long range charge-ordering has been observed to be associated with a phase transition between orthorhombic and monoclinic crystal systems.^{10,11} A schematic representation of the monoclinic structure of ilvaite is shown in Fig. 1. Edge-sharing double octahedra, $\text{FeO}_5(\text{OH})$, chains align parallel to the *c*-axis, with further Fe-centred octahedra sitting above and below the chain sharing four edges. Hydroxy-groups line the edges of the chains, hydrogen bonding to the apical sites of octahedra in neighboring chains. The Fe-containing chains are linked by disilicate, Si_2O_7 , units to form a three-dimensional structure containing seven-coordinated Ca^{2+} cations. In the charge-disordered

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† Electronic supplementary information (ESI) available: Tables of lattice parameters (Table S1) and BVS calculations (Table S2) used in the construction of Fig. 2 and 3. See DOI: 10.1039/c4tc00224e



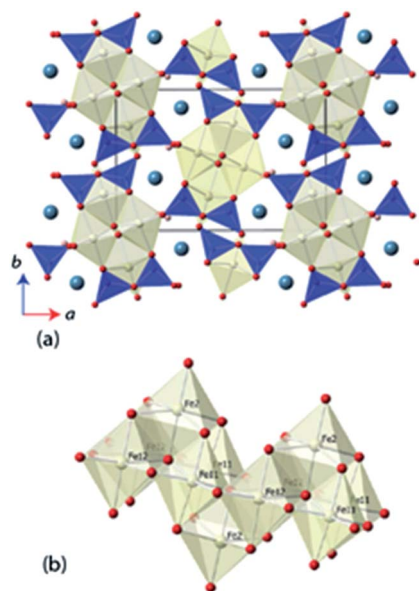


Fig. 1 (a) Structure of ilvaite viewed down the *c*-axis. Dark shaded tetrahedra are SiO_4 , and pale octahedra FeO_6 . Calcium is shown as a large sphere and hydrogen a small sphere. (b) The chains of edge-sharing Fe-centred octahedra showing the various iron sites as labelled.

orthorhombic form two distinct octahedral iron sites M1 and M2, exist where M1 is a mixed valence $\text{Fe}^{2+/3+}$ site, whilst M2 is entirely Fe^{2+} ordered; in natural ilvaites substitution of Mn^{2+} cations may occur into the M2 site.^{12,13} In the charge-ordered monoclinic form the M1 octahedral sites resolves as two independent crystallographic sites, M11 (Fe^{2+}) and M12 (Fe^{3+}).^{14,15} The degree of $\text{Fe}^{2+}/\text{Fe}^{3+}$ ordering within ilvaite also dictates the extent of deviation from orthorhombic symmetry, with β -angles up to 90.45° reported for complete ordering and $\beta = 90.00^\circ$ for entirely disordered $\text{Fe}^{2+}/\text{Fe}^{3+}$ sites.^{16,17} Ghose *et al.* investigated ilvaite below room temperature (305–5 K) using neutron powder diffraction (NPD) and showed charge-ordering between the M11 and M12 sites existed over this temperature range.¹⁸ A subsequent, higher temperature study using single crystal X-ray diffraction (295–400 K) extracted structural parameters at five 20 K intervals and above the transition temperature (T_{CO}) of ~ 346 K, changes in the Fe–O bond lengths within the Fe-octahedra, indicate complete charge disorder with the iron valence on both sites having converged towards 2.5+. Previous studies into the charge-ordering effects on structure in ilvaite and other materials, such as YBaCo_2O_5 , as a function of temperature have required a long data collection time at each temperature; this is particularly so for NPD data. This generally results in a coarse data set with structural information extracted at only a few temperatures around T_{CO} or reduced bond length accuracy from the use of X-ray diffraction methods. The lack of temperature resolution with data collected at only a few temperatures, either with single crystal X-ray diffraction or NPD, can be limiting in the investigation of the behavior of structures near T_{CO} . We have used the high-flux neutron powder diffractometer, D20 at the ILL, to extract rapidly, and with accurate structural resolution, a

full model of the long range charge-ordering and so determine the resulting structural changes in ilvaite below and through its monoclinic to orthorhombic transition temperature.

A sample of ilvaite (1.5 g, phase pure) from Dalnegorsk, Russia, was ground to a fine homogenous powder was studied on the D20 high flux neutron powder diffractometer at the ILL, Grenoble ($\lambda = 1.87 \text{ \AA}$).¹⁹ Following mounting in a cylindrical vanadium can, the sample was cooled to 140 K in a cryostat, allowing for a stabilisation time of 10 minutes, before ramping at 0.7 K min^{-1} to 380 K with a total data collection time of under 6 hours. With this relatively slow heating rate, little or no thermal gradient exists across the 1.5 g sample. Data were binned as 10 minute sets before further analysis; the temperature associated with each of these binned data sets was taken as that at mid-point of the data collection corrected for thermal lag between the measurement thermocouple and the sample.²⁰ To confirm the accuracy of this method very high quality NPD patterns were also collected on the D2B high resolution diffractometer ($\lambda = 1.594 \text{ \AA}$) at the Institut Laue Langevin, Grenoble, France.²¹ Diffraction patterns were collected over 180 minutes at four temperatures, 100, 298, 575 and 800 K, with samples mounted in cylindrical vanadium cans before being placed in a DISPLEX cryofurnace (100 K), suspended in the beam (298 K) or placed in a furnace (575, 800 K) before cooling/heating to the desired temperature. A thermal equilibration time of 15 minutes was allowed between each temperature change and data collection commencement. The unit cell parameters and atomic coordinates of ilvaite were extracted from each dataset through Rietveld profile refinement,²² using the GSAS suite of programs and the EXPGUI graphical user interface.^{23,24} The initial structural refinement against the data obtained at 100 K on D2B started with the monoclinic structure model described by Finger and Hazen.¹⁵ All atomic positions and ADPs were refined, and the M2 position was found to contain 30(1)% Mn. The final model obtained with the 100 K data was used as an initial model for the structural refinement of all other D2B data sets using the final model from the previous temperature block as a starting point for the next. The 575 and 800 K D2B data sets were refined against an orthorhombic initial model,²⁵ as the sample temperature was greater than T_{CO} . Similarly, the monoclinic 100 K – D2B model was used as an initial model for the first data set collected on D20, before carrying out a sequential GSAS refinement against all D20 data sets and extracting the required structural information.

Fig. 2 shows the evolution of the lattice parameters in ilvaite between 100 and 800 K focusing on the D20 data from 100 K to just above T_{CO} . Note that there is a very small systematic error between lattice parameter data extracted from D20 and D2B, likely to be due to instrumental differences, but overall values are in excellent agreement and trends in data readily resolved. The 7 K temperature resolution in the D20 data, derived from the data binning, coupled with the inherent precision in lattice parameters obtained from NPD (due to the strong scattering at high diffraction angles) allows the cell parameter behaviour as a function of temperature on approaching and exceeding a charge-ordering process to be



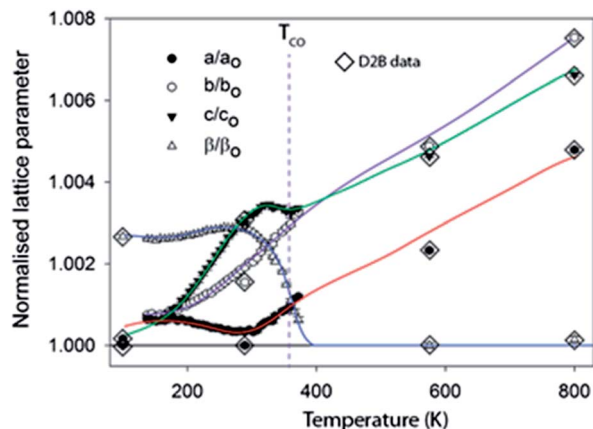


Fig. 2 Variation of the normalized, against the minimum value obtained, lattice parameters as a function of temperature.† Lines are best fit to the D20 data and have been extended as a guide to the eye into the regions covered by the D2B data.

deciphered much more clearly than previously.²⁵ As reported in this earlier work one of the most significant changes is in β , which decreases from a refined value in the monoclinic structure description of $90.252(5)^\circ$ (at its maximum at 250 K) to $90.05(2)^\circ$ as the cell becomes orthorhombic at T_{CO} . T_{CO} is estimated to be 352 K in this sample with 30% Mn on the M2 site close to the value reported by Ghose of 346 K for a sample with less than 1% Mn. However, between 150 K and 280 K β is observed to increase marginally as the structure expands rapidly in the c lattice direction but, notably, it contracts along a . Between 280 K and $T_{CO} = 352$ K as β decreases rapidly, the a -axis expands and the expansion along the c direction halts; the b lattice parameter expands normally over the full temperature range studied. Some faint evidence of these more subtle changes in lattice parameters can be seen in the work of Ghose;^{14,25} for example, an increase in a may be discerned near T_{CO} but the data were of insufficient precision and number to be confident of such trends. The anomalous behaviours of a , c and β near T_{CO} can be associated with the onset of charge ordering and these, plus the contraction along a and the very rapid expansion along c between 150 and 280 K, require further analysis of the bond lengths and derived charge distribution, *vide infra*. M–O distances were extracted from the crystallographic model refined against NPD data at each temperature. It is also worth noting that with NPD data the hydrogen position and the strength of the moderately strong O–H \cdots O hydrogen bond (O \cdots O \sim 2.6 Å) can also be monitored; however, the O–H bond length was found to behave normally over the full temperature range studied showing a small monotonic increase in length, *i.e.* no change in this part of the structure occurred on charge ordering. The derived M–O distances can be used to calculate bond valence sums (BVSs) for individual metal sites according to the method of Brown and Altermatt²⁶ and, therefore, assign charges each site. For these calculations r_0 values for both Fe^{3+} and Fe^{2+} were initially explored for each unique iron site and the BVSs obtained used to identify the associated valence

state. Thus the use of the Fe^{2+} r_0 value for a site occupied by Fe^{3+} leads to BVS values around 2.3 to 2.4 + instead of the expected 2.0 ± 0.15 (if occupied by Fe^{2+}), indicating that the Fe^{3+} r_0 value is applicable for this site. Similarly the use of the r_0 value for Fe^{3+} for a site in reality occupied by Fe^{2+} also yields values near 2.4+ rather than the expected 2.0 ± 0.15 expected with the correct choice of r_0 for Fe^{2+} . Note that the use of the incorrect r_0 value is clearly shown in these calculations but the BVSs obtained with the correct r_0 , while lying close to the expected BVS ($2+$ or $3+$), do not fully reach the value of the higher or lower oxidation state BVS. This indicates some residual structural strain in the material after charge-ordering occurs in that the oxide ion geometry cannot fully relax, so an Fe^{2+} site remains compressed and the Fe^{3+} site coordination over-expanded. Furthermore BVSs are expected to decrease slightly with increasing temperature due to thermal expansion.

For the monoclinic form at 100 K the BVS analysis clearly assigned M2 (70% Fe, 30% Mn) as purely M^{2+} and M11 as Fe^{2+} and M12 as Fe^{3+} in agreement with the previous reports.²⁵ Calculations of BVSs were then undertaken with appropriate r_0 values for all 40 temperatures and Fig. 3 presents these data for the three independent metal sites of the monoclinic phase. For the orthorhombic phase above T_{CO} M11 and M12 become equivalent and BVS were carried out for the single site using Fe^{2+} and Fe^{3+} r_0 values and also using an r_0 value estimated for $Fe^{2.5+}$; all three BVS calculated values are plotted in Fig. 3. The charge ordering in ilvaite is clearly seen in this plot and occurs gradually between 300 and 380 K and is essentially complete at $T_{CO} = 362$ K in agreement with the lattice parameter data. Below T_{CO} M11 is occupied by Fe^{3+} while the electron charge localises on M12 producing Fe^{2+} . In both monoclinic and orthorhombic forms the BVSs decrease slightly with the increasing temperature as expected due to thermal expansion and increases in Fe–O distances.

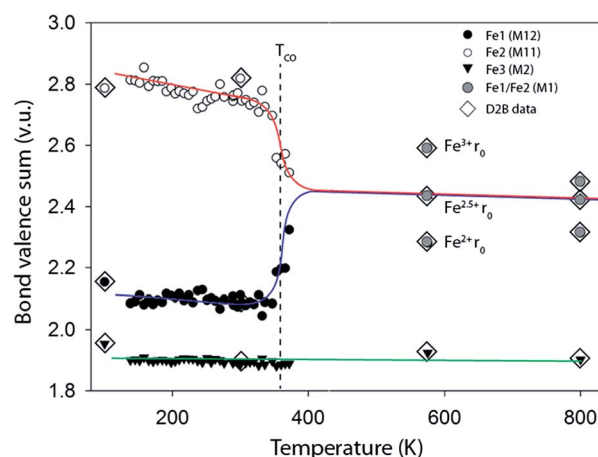


Fig. 3 Bond valence sums for the three transition metal sites in ilvaite calculated using the Fe^{3+} r_0 values for M11 and Fe^{2+} r_0 parameter for M12 and M2.† For the two high temperature data points in the orthorhombic phase where M11 is equivalent to M12, M1, with an average oxidation state of 2.5+ the BVS was calculated with Fe^{2+} , $Fe^{2.5+}$ and Fe^{3+} r_0 values.



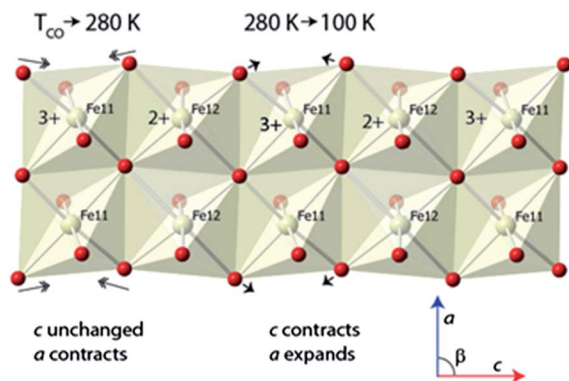


Fig. 4 Ordered Fe11 and Fe12 layer in the charge ordered state of ilvaite indicating the changes that occur in the atomic positions and lattice parameters on cooling from 280 to 150 K Fe^{2+} .

The anomalous behaviours in the lattice parameters noted earlier can now also be interpreted. Considering the behaviour as a sample of ilvaite is cooled from 360 and 280 K and the charge-ordering occurs and the a and b axes contract normally while c is almost invariant; this behaviour is associated with the rapid relocation of the edge-shared oxygen atoms along c and closer to M11 with its higher formal charge Fe^{3+} and away from M12, Fig. 4. This relocation is, as would be expected with equal numbers of bond lengthening and shortening, almost distance neutral along c , while a and b contract with decreasing thermal motion. The BVS for the M11 site increases most rapidly on cooling between 280 K and 100 K indicating that over this temperature range the loss of thermal energy after charge-ordering has occurred causes the oxide anions to continue to move towards this site while the BVs for M12 indicate that distances around this site change little on further cooling involving a combination of a slight decrease along c but increase along a . This further significant displacement towards M11 but almost unchanged M12–O distances causes a “concertina” effect and is responsible for the observed changes in lattice parameters seen Fig. 1. The c -lattice parameter contracts rapidly as a result of continued reduction in M11–O distances along this direction while a expands as the oxygen maintains its O–M12 distance. Fig. 4 shows a portion of the double chain of edge sharing octahedra centred on M11 and M12 in the ilvaite structure and the major displacements in oxide ion positions that occur at just below T_{CO} and then further cooling from 280 K to 150 K. These displacements lie mainly in the ac plane while the b direction remains almost unaffected and contracts normally on cooling.

Conclusions

In summary NPD data collected over a 6 hour period on a high flux instrument on ilvaite allows the electron charge-ordering to be directly probed and the structural ramifications of this process below T_{CO} to be investigated in detail. This information provides unprecedented insights into the charge-ordering process to be obtained, not only close to the charge-ordering

temperature but also in the charge-ordered state across a wide temperature range. The subtle changes in structure that are found in the charge-ordered ilvaite, even at 100–200 K below T_{CO} , indicate that structural relaxation occurs well before or after the electron delocalization on heating or cooling respectively.

Consideration of published investigations of other iron compounds that demonstrate charge ordering^{2–5,7} show that, in general, the resolution in terms of the number of temperatures studied and structural accuracy and precision obtained is much lower than with very high flux NPD analysis used in this work. Thus small variations in lattice parameters and bond lengths that occur above and below T_{CO} have not been reported previously. For example, Fe_2OBO_3 ,⁴ where only two temperatures were studied. One exception is with studies of $\text{LnBaFe}_2\text{O}_5$ ($\text{Ln} = \text{Ho}, \text{Nd}$) phases,² where subtle changes in lattice parameters above T_{CO} were observed, using the D20 instrument, and explained in terms of a premonitory transition; however, in this case full profile refinement and structural parameter extraction across the transition was not undertaken. The application of the very high-flux, variable temperature (with a small temperature step) NPD technique would be applicable to the study of many other materials that undergo charge-ordering and produce a much better understanding of this phenomenon and structural relaxation that might occur above and below T_{CO} .

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Notes and references

- 1 J. Rodriguez-Carvajal, M. Hennion, F. Moussa, A. H. Moudeen, L. Pinsard and A. Revcolevschi, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1998, **57**, R3189; F. Damay, C. Martin, M. Hervieu, A. Maignan, B. Raveau, G. André and F. Bourée, *J. Magn. Magn. Mater.*, 1998, **184**, 71; J. F. Mitchell, D. N. Argyriou, J. D. Jorgensen, D. G. Hinks, C. D. Potter and S. D. Bader, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1997, **5**, 63.
- 2 P. M. Woodward, D. E. Cox, E. Moshopoulou, A. W. Sleight and S. Morimoto, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2000, **62**, 844.
- 3 D. K. Pratt, S. Chang, W. Tian, A. A. Taskin, Y. Ando, J. L. Zarestky, A. Kreyssig, A. I. Goldman and R. J. McQueeney, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2013, **87**, 045127.
- 4 M. Angst, P. Khalifah, R. P. Hermann, H. J. Xiang, M.-H. Whangbo, V. Varadarajan, J. W. Brill, B. C. Sales and D. Mandrus, *Phys. Rev. Lett.*, 2007, **99**, 086403.
- 5 N. Ikeda, *et al.*, *Nature*, 2005, **436**, 1136.
- 6 K. M. Shen, *et al.*, *Science*, 2005, **307**, 901.
- 7 J. P. Wright, J. P. Attfield and P. G. Radaelli, *Phys. Rev. Lett.*, 2001, **87**, 266401.



- 8 A. M. M. Abeykoon, E. S. Božin, W.-G. Yin, G. Gu, J. P. Hill, J. M. Tranquada and S. J. L. Billinge, *Phys. Rev. Lett.*, 2013, **111**, 096404.
- 9 T. Vogt, P. M. Woodward, P. Karen, B. A. Hunter, P. Henning and A. R. Moodenbaugh, *Phys. Rev. Lett.*, 2000, **84**, 2969.
- 10 N. V. Belov and V. I. Mokeeva, *Trudy Inst. Krist. Akad. Nauk SSSR*, 1954, **9**, 89.
- 11 P. Bartholome, J. C. Duchesne and L. Van der Plas, *Ann. Soc. Geol. Belg.*, 1968, **90**, 779.
- 12 N. Haga and Y. Takeuchi, *Z. Krist.*, 1976, **144**, 161.
- 13 P. Bonazzi and L. Bindi, *Am. Mineral.*, 2002, **87**, 845.
- 14 D. A. Nolet and R. G. Burns, *Phys. Chem. Miner.*, 1979, **4**, 221.
- 15 L. W. Finger and R. M. Hazen, *Z. Krist.*, 1987, **179**, 415.
- 16 Y. Takeuchi, N. Haga and M. Bunno, *Z. Krist.*, 1983, **163**, 267.
- 17 Y. Takeuchi, H. Sawada and H. Taniguchi, *Proc. Inst. Nat. Sci., Nihon Univ.*, 1993, **28**, 39.
- 18 S. Ghose, A. W. Hewat and M. Marezio, *Phys. Chem. Miner.*, 1984, **11**, 67.
- 19 T. C. Hansen, P. F. Henry, H. E. Fischer, J. Torregrossa and P. Convert, *Meas. Sci. Technol.*, 2008, **19**, 034001.
- 20 M. T. Weller, P. F. Henry, V. P. Ting and C. C. Wilson, *Chem. Commun.*, 2009, 2973.
- 21 A. W. Hewat and S. Heathman, *Acta Crystallogr., Sect. A: Found. Crystallogr.*, 1984, **40**, C364.
- 22 H. M. Rietveld, *J. Appl. Crystallogr.*, 1969, **2**, 65.
- 23 B. H. Toby, *J. Appl. Crystallogr.*, 2001, **34**, 210.
- 24 A. C. Larson and R. B. Von Dreele, *General Structure Analysis System (GSAS)*, Los Alamos National Laboratory Report LAUR, 1994, pp. 86–748.
- 25 S. Ghose, K. Tsukimura and D. M. Hatch, *Phys. Chem. Miner.*, 1989, **16**, 483.
- 26 I. D. Brown and D. Altermatt, *Acta Crystallogr., Sect. B: Struct. Sci.*, 1985, **41**, 244.

